

**10B.3      DIAGNOSING THE INTERCEPT PARAMETER FOR EXPONENTIAL RAINDROP SIZE DISTRIBUTION BASED ON VIDEO DISDRMETER OBSERVATIONS**

Guifu Zhang<sup>1</sup>, Ming Xue<sup>1,2</sup>, Qing Cao<sup>1</sup>, and Daniel Dawson<sup>1,2</sup>  
<sup>1</sup>School of Meteorology and <sup>2</sup>Center for Analysis and Prediction of Storms  
 University of Oklahoma, Norman OK 73072

**1. Introduction**

Information about the drop size distribution (DSD) is essential for understanding precipitation physics, estimating rainfall, and improving microphysics parameterizations in numerical weather prediction (NWP) models (Steiner et al. 2004). The characteristics of rain DSDs are often associated with the types of storms (e.g., convective versus stratiform rain) and their stages of development (e.g., the developing versus decaying stage, Brandes et al. 2006). Strong convective rain usually contains both large and small drops and has a broad DSD while the decaying stage is often dominated by small drops. Stratiform rain usually contains relatively larger drops but has a low number concentration for a given rain rate (Zhang et al. 2006).

Rain DSDs are usually represented by distribution models such as the exponential distribution, gamma distribution, and lognormal distribution. A DSD model usually contains a few free parameters that should be easy to determine and the model should be capable of capturing the main physical processes and properties. The exponential distribution is the most commonly used DSD model that has some of these properties, and it is given by

$$N(D) = N_0 \exp(-\Lambda D). \quad (1)$$

It contains two free parameters,  $N_0$  and  $\Lambda$ . A single-moment bulk microphysics model predicts one of the moments of the DSD which determines one of the two parameters. The intercept parameter  $N_0$  is usually specified so that  $\Lambda$  is uniquely related to the water content,  $W$ , which is the 3<sup>rd</sup> moment of the DSD that is predicted. The Marshall–Palmer (M-P, Marshall and Palmer 1948) exponential DSD model with the  $N_0$  value fixed at  $8000 \text{ m}^{-3} \text{ mm}^{-1} = 8 \times 10^6 \text{ m}^{-4}$  is widely used for representing warm rain (Kessler 1969) as well as ice (e.g., Lin et al. 1983; Hong et al. 2004) microphysics.

However, disdrometer observations and numerical model simulations indicate that  $N_0$  and number concentration ( $N_t$ ) are not constant, but vary depending on pre-

cipitation type, rain intensity and stage of development. Sauvageot and Lacaux (1995) showed variations of both  $N_0$  and  $\Lambda$  from impact disdrometer measurements. Recent observations by 2D Video Disdrometers (2DVD) suggest that rain DSDs are better represented by a constrained Gamma distribution (Zhang et al. 2001) that also contains two free parameters. In Zhang et al. (2006), the constrained Gamma model was further simplified to a single parameter model for bulk microphysical parameterization and the model produced more accurate precipitation system forecasts than the M-P model. Since the exponential distribution model is widely used, a diagnostic relation of  $N_0$  as a function of  $W$  would improve rain estimation and microphysical parameterization that are based on such an improved model.

In this study, we derive a diagnostic  $N_0$  relation from rain DSD data that were collected in Oklahoma using disdrometers. To minimize the error effects introduced in the fitting procedure, we formulate the problem with a relation between two DSD moments. A diagnostic relation is found from the relation between two middle moments. Section 2 describes methods of deriving the diagnostic relation and section 3 presents results of diagnosing  $N_0$  from water content using 2DVD measurements. In Section 4, we discuss applications of the diagnostic relation in the parameterization of rain physics and microphysical processes. Final summary and discussions are given in section 5.

**2. Diagnosing methods**

The diagnostic relation for the intercept parameter  $N_0$  as a function of water content can be derived using two different approaches.

One approach is to first find the DSD parameters ( $N_0$ ,  $\Lambda$ ) by fitting DSD (e.g., disdrometer) data to the exponential function (1) for each DSD, and then plot the estimated  $N_0$  versus  $W$  for the whole dataset for fitting a mean relation.

The  $n^{\text{th}}$  moment of the exponential DSD (1) is

$$M_n = \int D^n N(D) dD = N_0 \Lambda^{-(n+1)} \Gamma(n+1) \quad (2)$$

Hence, the DSD parameters,  $N_0$  and  $\Lambda$ , can be determined from any two moments ( $M_l$ ,  $M_m$ ) as

---

<sup>1</sup>Corresponding author: Dr. Guifu Zhang  
 School of Meteorology, University of Oklahoma,  
 120 David L. Boren Blvd, Suite 5900, Norman OK 73072  
 E-mail: guzhang1@ou.edu

$$\Lambda = \left( \frac{M_l \Gamma(m+1)}{M_m \Gamma(l+1)} \right)^{\frac{1}{m-l}}, \quad (3)$$

$$N_0 = \frac{M_l \Lambda^{l+1}}{\Gamma(l+1)}. \quad (4)$$

When  $N_0$  is obtained along with the water content  $W$  for DSD data sets, a  $N_0 - W$  relation can be found through another fitting procedure, e.g., the power-law fitting. It is noted that the values of the estimated  $N_0$  depend on which two moments are used and on the accuracy of the two moment estimates. Since the estimates of both the moments ( $M_l, M_m$ ) have error, the DSD parameters ( $N_0, \Lambda$ ) obtained from them will also have error. The natural variation in DSDs also causes a large scatter in the  $N_0 - W$  plot (see Fig.1 in next section). Hence, the  $N_0 - W$  relation derived from the above procedure tends to have larger errors.

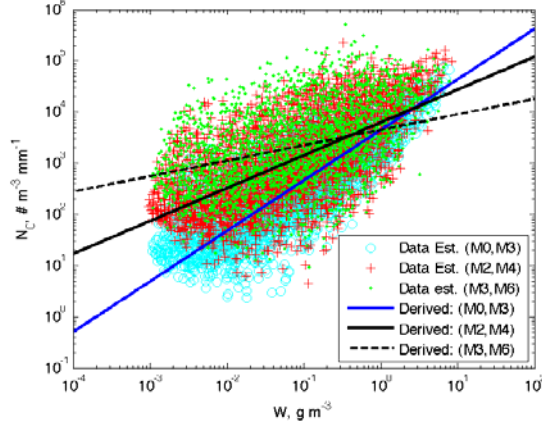


Fig. 1: Dependence of intercept parameter ( $N_0$ ) on water content ( $W$ ). Scattered points are fitted results from a pair of DSD moments. Straight lines are derived relations using the moment relation method.

To minimize the error effects introduced in the fitting procedures, we propose an alternative procedure for obtaining the  $N_0 - W$  relation. Here, we first seek to establish a relation between two DSD moments. With this relation, the exponential distribution is reduced to having a single free parameter so that  $N_0$  can be determined from  $W$ . Suppose that two DSD moments  $M_l$  and  $M_m$  are related by a power-law relation:

$$M_l = aM_m^b, \quad (5)$$

where  $a$  and  $b$  are coefficients that can be estimated from disdrometer observations.

From Eq. (2) for the third moment, we have water content  $W = \frac{\pi}{6} \rho M_3 = \pi \rho N_0 \Lambda^{-4}$ , [ $\rho$  is water density],

yielding the slope parameter  $\Lambda = \left( \frac{N_0 \pi \rho}{W} \right)^{1/4}$ . Substituting (2) into (5) for  $M_l$  and  $M_m$ , and making use of the relation for  $\Lambda$ , we obtain

$$N_0 = \alpha W^\beta, \quad (6)$$

where

$$\alpha = \left( a \frac{\Gamma^b(m+1)}{\Gamma(l+1) \pi^c \rho^c} \right)^{\frac{1}{1-b+c}}, \quad (7)$$

$$\beta = \frac{c}{1-b+c}, \quad (8)$$

and

$$c = \frac{b(m+1) - (l+1)}{4}. \quad (9)$$

Hence, (6) - (9) constitute a general formulation for deriving a  $N_0 - W$  relation using a statistical relation between two DSD moments. When the coefficients  $a$  and  $b$  in the relation (5) are determined from a set of DSD data, we have a diagnostic relation between the water content  $W$  and the intercept parameter  $N_0$ . This is the procedure that will be used in the next section with a disdrometer data set.

### 3. Derivation of the $N_0 - W$ relation from disdrometer observations

We test our method for deriving the  $N_0 - W$  relation using disdrometer data collected in Oklahoma during the summer seasons of 2005 and 2006 (Cao et al. 2007). Three 2DVDs, operated respectively by the University of Oklahoma (OU), National Center for Atmospheric Research (NCAR) and National Severe Storms Laboratory (NSSL) were deployed at the NSSL site in Norman, Oklahoma, and at the Southern Great Plains (SGP) site of the Atmospheric Radiation Measurement (ARM) program. The ARM site is located approximately 28 km south of the NSSL site. A total of 58 days of disdrometer data covering a total of 13379 minutes of rainfall periods with total drop counts greater than 10 were collected. Among them are 1092 minutes of data that have side-by-side measurements by two 2DVDs. The recorded raindrops within each minute were processed to produce one-minute DSD samples, yielding a total of 13379 DSDs.

With the side-by-side data, measurement errors of DSDs were quantified. The sampling errors are further reduced by sorting and averaging based on two parameters (SATP), a method that combines DSDs with similar rainfall rates ( $R$ ) and median volume diameters ( $D_0$ ) (Cao et al. 2007). There are 1092 quality-controlled DSDs after SATP processing for the dataset. The DSD moments are estimated by the sum of weighted DSDs as defined in (2). Cao et al. (2007) showed that the middle

moments ( $M_2$ ,  $M_3$  and  $M_4$ ) have the least error (see their Table 1).

The exponential DSD parameters  $N_0$  and  $\Lambda$  are estimated from the moment pairs of ( $M_0$ ,  $M_3$ ), ( $M_2$ ,  $M_4$ ), and ( $M_3$ ,  $M_6$ ) using Eqs. (3) and (4). The exponentially-fitted  $N_0$  values are plotted versus the rain water content in Fig. 1. As expected, there is a large scatter in the  $N_0 - W$  plot because of measurement and model errors as well as natural variations. Also, different moment pairs produce different results of  $N_0$  due to differences in estimation error and error propagation in the fitting procedure. It is clear that there is a positive correlation between  $N_0$  and  $W$ . However, there will be high uncertainty if we were to fit the data points of this plot directly to a  $N_0 - W$  relation.

Instead, an  $N_0 - W$  relation is derived from a moment relation as outlined in section 2. Figure 2 shows the scatter plots of moment estimates for the pair of ( $M_2$ ,  $M_4$ ), and a power-law relation is obtained as

$$M_2 = 1.42M_4^{0.836}, \quad (10)$$

i.e.,  $a = 1.42$  and  $b = 0.836$  in Eq.(5). The correlation between the moments is high, with a correlation coefficient of 0.88 in the linear domain and 0.93 in the logarithmic domain. Substituting for  $a$  and  $b$  in (6)-(9), we obtain  $\alpha = 6257$  and  $\beta = 0.642$ , therefore

$$N_0(M_2, M_4) = 6257W^{0.642}. \quad (11)$$

This  $N_0 - W$  relation is shown in Fig. 1 along with those derived from moment pairs ( $M_0$ ,  $M_3$ ) and ( $M_3$ ,  $M_6$ ). The lower (higher) moment pair yields a relation with a larger (smaller) power. As discussed earlier, the middle moment pair ( $M_2$ ,  $M_4$ ) has smaller error and therefore relation (11) is recommended.

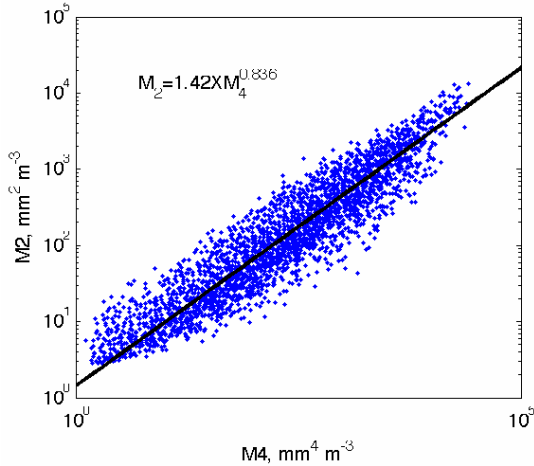


Fig. 2: Inter-relationships among DSD moments based on disdrometer measurements. Scattered points are direct estimates from disdrometer measurements. Straight lines represent fitted power-law relations.

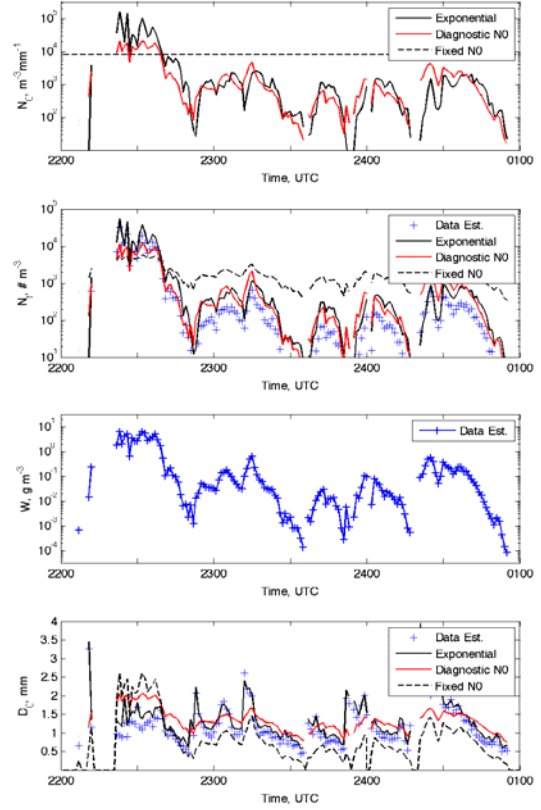


Fig. 3: Time series comparison of physical parameters: intercept parameter ( $N_0$ ), total number concentration ( $N_1$ ), water content ( $W$ ), and median volume diameter ( $D_0$ ) for a rain event starting on July 21, 2006. Results are shown for disdrometer measurements and fitted values using exponential, diagnostic- $N_0$ , and fixed- $N_0$  DSD models.

For a better understanding of the  $N_0 - W$  relation (11), Figure 3 shows an example of  $N_0$  values along with other physical parameters ( $N_1$ ,  $W$ , and  $D_0$ ) as a function of time for a rain event starting on July 21, 2006. It was a strong convective storm followed by stratiform rain as well as weak convection passing over the OU disdrometer deployed at the ARM site at Washington, Oklahoma. The water content is very low during the stratiform precipitation period, but the median volume diameter  $D_0$  is comparable to that of strong convection. The comparison between exponentially-fitted  $N_0$  values from DSD moments  $M_2$  and  $M_4$  and those diagnosed from  $W$  using (11) is plotted in Fig. 3a. As shown in Fig. 3b, the moment fitting of the exponential DSD model yields a good estimate of total number concentration  $N_1$  as compared with the direct estimates from DSD data (discrete “+”). Here, the fitted  $N_0$  can be considered as “truth” because  $N_0$  is a model parameter which is obtained through the fitting procedure of Eqs. (2) - (4). It is clear that the diagnosed  $N_0$  captures the main trend of the observed rain storm very well in a dynamic range

of more than two orders of magnitude; that is, from an order of  $10^4$  for strong convection to 10 for light stratiform precipitation. In comparison, the fixed- $N_0$  M-P model overestimates  $N_0$  except for heavy convective rain. Figure 3d compares median volume diameter  $D_0$  calculated from the DSD data, estimated using the diagnostic- $N_0$  model and that with the fixed- $N_0$  model. Again, the diagnostic  $N_0 - W$  relation produces a much better agreement with the measurements than does the fixed- $N_0$  model.

#### 4. Application to warm rain microphysical parameterization

The warm rain microphysical processes related to the DSD include rain evaporation, accretion of cloud water by rain water, and rain sedimentation. The microphysical parameters of these processes based on the exponential DSD model have been derived by Kessler (1969: Table 4). After unit conversion, the evaporation rate ( $R_e$  in  $\text{kg kg}^{-1} \text{s}^{-1}$ ), accretion rate ( $R_c$  in  $\text{kg kg}^{-1} \text{s}^{-1}$ ), mass-weighted terminal velocity ( $V_{tm}$  in  $\text{m s}^{-1}$ ), and reflectivity factor ( $Z$  in  $\text{mm}^6 \text{m}^{-3}$ ) are obtained as

$$R_e = 2.17 \times 10^{-5} E_e N_0^{7/20} (q_{vs} - q_v) W^{13/20}, \quad (12a)$$

$$R_c = 1.65 \times 10^{-3} E_c N_0^{1/8} q_c W^{7/8}, \quad (12b)$$

$$V_{tm} = 16.4 N_0^{-1/8} W^{1/8} (\rho_0 / \rho)^{0.5}, \quad (12c)$$

$$Z = 1.73 \times 10^7 N_0^{-3/4} W^{7/4}. \quad (12d)$$

where  $E_e$  and  $E_c$  are the evaporation and accretion efficiency factors, respectively (normally taken as 1),  $W$  is rain water content in  $\text{g m}^{-3}$  as before ( $W = 1000 \rho q_r$ ),  $q_v$ ,  $q_c$  and  $q_r$  are, respectively, the water vapor, cloud water and rain water mixing ratios in  $\text{kg kg}^{-1}$ .

Substituting the diagnostic relation (11) into (12) and assuming unit saturation deficit and unit cloud water mixing ratio as well as unit efficiency factors, we obtain a parameterization scheme based on the diagnostic  $N_0$ . The terms corresponding to those in Eq. (12) are listed in Table 1 along with those of the standard fixed- $N_0$  M-P model. The coefficients of these terms are similar for the two schemes, but the powers are substantially different. The larger power in evaporation rate means more (less) evaporation for heavy (light) rain compared to the fixed- $N_0$  model. The smaller power in the reflectivity formula for the diagnostic- $N_0$  model gives smaller (larger) reflectivity than the fixed  $N_0$  for heavy (light) rain. This may lead to a better agreement between numerical model forecasts and radar observations. The former tends to over-predict large reflectivity values and under-predict low reflectivity values. In this sense, the diagnostic- $N_0$  model has similar properties as the simplified-constrained-gamma model investigated in Zhang et al. (2006)

Table 1: Parameterization of warm rain processes with diagnostic  $N_0$  and fixed  $N_0$

Parameterized quantity	Diagnostic $N_0$	Fixed $N_0$
$N_0, \text{m}^{-3} \text{mm}^{-1}$	$6257W^{0.642}$	8000
$R_e, \text{kg kg}^{-1} \text{s}^{-1}$	$4.63 \times 10^{-4} W^{0.875}$	$5.03 \times 10^{-4} W^{0.65}$
$R_c, \text{kg kg}^{-1} \text{s}^{-1}$	$4.92 \times 10^{-3} W^{0.955}$	$5.08 \times 10^{-3} W^{0.875}$
$V_{tm}, \text{m s}^{-1}$	$5.50W^{0.0447}$	$5.32W^{0.125}$
$Z, \text{mm}^6 \text{m}^{-3}$	$2.46 \times 10^4 W^{1.27}$	$2.04 \times 10^4 W^{1.75}$

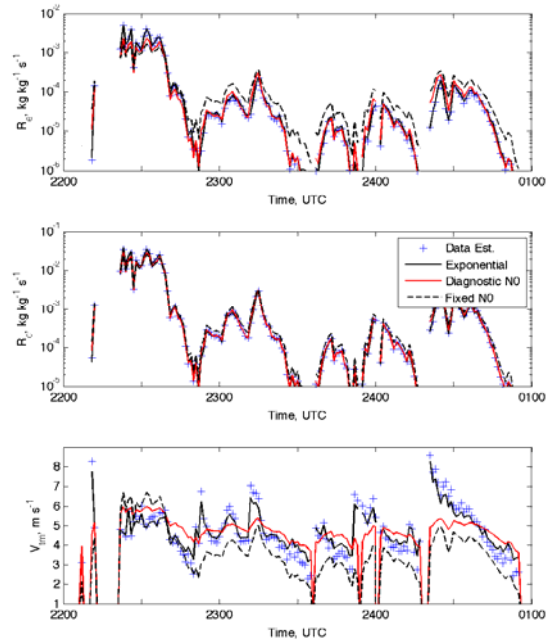


Fig. 4: As in Fig. 3 except for evaporation rate for a unit vapor saturation deficit ( $R_e$ ), accretion rate ( $R_c$ ) for a unit cloud water content, and mass-weighted terminal velocity ( $V_{tm}$ ) DSD models.

Figure 4 compares the terms for the microphysical processes estimated from  $W$  using the diagnostic- $N_0$  DSD model with those from the fixed- $N_0$  M-P DSD model. Direct calculations from the observed DSD data and those fitted with the exponential model with  $N_0$  as one of the two free parameters are also shown for reference. The results may appear to be close to each other in the semi-logarithm plots but actually, the fixed- $N_0$  model overestimates the evaporation rate for stratiform rain by about a factor of five and underestimates that for strong convection. This might be the reason that the parameterization coefficients in the Kessler scheme are sometimes reduced by a half and more in order to obtain

a better match of modeling results with observations (e.g., Sun and Crook 1997). The diagnostic- $N_0$  model therefore characterizes rain evaporation, accretion and rainfall processes more accurately than the fixed- $N_0$  model for both heavy and light rainfall. This is because by introducing the dependency of  $N_0$  on  $W$  based on observations, raindrop number concentration and total surface area of rain drops are better represented, leading to a better estimation of evaporation and accretion rates.

## 5. Summary and Discussions

In this paper, we present a way of diagnosing the intercept parameter  $N_0$  of the exponential distribution and apply the diagnostic- $N_0$  DSD model to improving warm rain microphysical parameterization. The diagnostic relation is derived from a relation between two DSD moments that are estimated from 2D video disdrometer data. The DSD data were collected in Oklahoma during the summer seasons of 2005 and 2006, which should be representative for rains in the central Great Plains region. The  $N_0 - W$  relation is used to improve the Kessler parameterization scheme of warm rain microphysics, and can be used in schemes containing ice-phases also (e.g., Lin et al. 1983; Hong et al. 2004).

It has been shown that the diagnostic- $N_0$  model better characterizes natural rain DSDs, including the physical properties (e.g.,  $N_t$  and  $D_0$ ) and microphysical processes. For a given water content, the diagnostic- $N_0$  DSD model represents the total number concentration, median volume diameter, reflectivity factor, evaporation rate and accretion rate much more accurately than the M - P model with a fixed  $N_0$ . Compared with the M-P model-based Kessler scheme, the modified parameterization scheme with a diagnostic  $N_0$  has the following advantages: (i) it leads to less (more) evaporation for light (heavy) rain and therefore can preserve stratiform rain better in numerical models, and (ii) it yields a larger (smaller) reflectivity factor for light (heavy) rain, having the potential of yielding a better agreement between model predicted and radar observed reflectivities in a similar way as the simplified constrained gamma model. Realistic simulation of reflectivity is important for assimilating radar reflectivity data into NWP models.

It is noted that the diagnostic  $N_0 - W$  relation obtained in this paper is based on a specific set of disdrometer data in the summer season of a specific climate region. While the methodology developed in this paper is general, the coefficients in the relation may need to be tuned to better fit specific regions and/or seasons or specific rain types. The improved parameterization based on the diagnostic- $N_0$  model is now being tested within the Advanced Regional Prediction System (ARPS, Xue

et al. 2003) to examine its impact on precipitation forecast; the results will be presented in a future paper.

*Acknowledgments* Authors greatly appreciate helps from Drs. Edward Brandes, Terry Schuur, Robert Palmer, Phillip Chilson and Ms. Kyoko Iketa. The sites for disdrometer deployment at the Kessler farm were provided by Atmospheric Radiation Measurement (ARM) Program. This work was primarily supported by a NSF grant ATM-0608168. Dan Dawson were also supported by NSF grants ATM-0530814 and Ming Xue further by ATM-0331594 and ATM-0331756.

## References

- Brandes, E. A., G. Zhang, and J. Sun, 2006: On the influence of assumed drop size distribution form on radar-retrieved thunderstorm microphysics. *J. Appl. Meteor.*, **45**, 259-268.
- Cao, Q., G. Zhang, E. Brandes, T. Schuur, A. Ryzhkov, and K. Ikeda, 2007: Analysis of video disdrometer and polarimetric radar data to characterize rain microphysics in Oklahoma. *J. Appl. Meteor. Climat. Under review.*
- Hong, S.-Y., J. Dudhia, and S.-H. Chen, 2004: A revised approach to ice microphysical processes for the bulk parameterization of clouds and precipitation. *Mon. Wea. Rev.*, **132**, 103-120.
- Kessler, E., 1969: On the distribution and continuity of water substance in atmospheric circulations. *Meteor. Monogr.*, No. 32, Amer. Meteor. Soc., 84 pp.
- Lin, Y.-L., R. D. Farley, and H. D. Orville, 1983: Bulk parameterization of the snow field in a cloud model. *J. Climate Appl. Meteor.*, **22**, 1065-1092.
- Marshall, J.S., and W. McK. Palmer, 1948: The distribution of raindrops with size, *J. Meteor.*, **5**, 165-166.
- Sauvageot, H., and J.-P. Lacaux, 1995: The shape of averaged drop size distributions. *J. Atmos. Sci.*, **52**, 1070-1083.
- Steiner, M., J. A. Smith, and R. Uijlenhoet, 2004: A microphysical interpretation of radar reflectivity-rain rate relationships. *J. Atmos. Sci.*, **61**, 1114-1131.
- Sun, J, and N. A. Crook, 1997: Dynamical and microphysical retrieval from Doppler radar observations using a cloud model and its adjoint. Part I: Model development and simulated data experiments. *J. Atmos. Sci.*, **54**, 1642-1661.
- Xue, M., D.-H. Wang, J.-D. Gao, K. Brewster, and K. K. Droegemeier, 2003: The Advanced Regional Prediction System (ARPS), storm-scale numerical weather prediction and data assimilation. *Meteor. Atmos. Physics*, **82**, 139-170.
- Zhang, G., J. Vivekanandan, and E. Brandes, 2001: A method for estimating rain rate and drop size distribution from polarimetric radar measurements. *IEEE Trans. Geosci. Remote Sens.*, **39**, 830-841.
- Zhang, G., J. Sun, and E. A. Brandes, 2006: Improving parameterization of rain microphysics with disdrometer and radar observations. *J. Atmos. Sci.*, **63**, 1273-1290.